

[7] U dvostranom modelu za slušaj $u_1 = -u_3$ i $k^2 = \lambda^2$, nađite ovisnost ψ_3' (ψ_1') za totalno stabilne valove ($\delta = 0$).

Pi: diferencijalne jednačine za 2LM:

$$1) \left(\frac{\partial}{\partial t} + u_m \frac{\partial}{\partial x} \right) \frac{\partial^2 \psi_m}{\partial x^2} + \beta \frac{\partial \psi_m}{\partial x} + u_T \frac{\partial}{\partial x} \left(\frac{\partial^2 \psi_T}{\partial x^2} \right) = 0$$

$$2) \left(\frac{\partial}{\partial t} + u_m \frac{\partial}{\partial x} \right) \left(\frac{\partial^2 \psi_T}{\partial x^2} - 2\lambda^2 \psi_T \right) + \beta \frac{\partial \psi_T}{\partial x} + u_T \frac{\partial}{\partial x} \left(\frac{\partial^2 \psi_m}{\partial x^2} + 2\lambda^2 \psi_m \right) = 0$$

- probna rješenja: $\psi_m = A e^{ik(x-ct)}$
 $\psi_T = B e^{ik(x-ct)}$ } \Rightarrow ubacimo u 1) i 2)

$$\Rightarrow (1) [(c - u_m)k^2 + \beta]A - u_T k^2 B = 0$$

$$(2) -u_T (k^2 - 2\lambda^2)A + [(c - u_m)(k^2 + 2\lambda^2) + \beta]B = 0$$

\Rightarrow dobije se disperzijska relacija:

$$c = u_m - \frac{\beta(k^2 + \lambda^2)}{k^2(k^2 + 2\lambda^2)} + \underbrace{\left[\frac{\beta^2 \lambda^4}{k^4(k^2 + 2\lambda^2)^2} - \frac{u_T^2(2\lambda^2 - k^2)}{k^2 + 2\lambda^2} \right]}_{\delta}$$

- zadržat kaže $u_1 = -u_3 \Rightarrow u_m = \frac{u_1 + u_3}{2} = 0$; $u_T = \frac{u_1 - u_3}{2} = u_1 = -u_3$

- znamo: $\psi_m = \frac{\psi_1' + \psi_3'}{2}$; $\psi_T = \frac{\psi_1' - \psi_3'}{2}$

$$\Rightarrow \frac{\psi_m}{A} = \frac{\psi_T}{B} \Rightarrow \frac{\psi_1' + \psi_3'}{2A} = \frac{\psi_1' - \psi_3'}{2B} \Rightarrow \psi_1' + \psi_3' = \frac{A}{B} \psi_1' - \frac{A}{B} \psi_3'$$

$$\Rightarrow \left(1 + \frac{A}{B}\right) \psi_3' = \left(\frac{A}{B} - 1\right) \psi_1' \Rightarrow \psi_3' = \frac{\frac{A}{B} - 1}{\frac{A}{B} + 1} \psi_1' \Rightarrow \boxed{\psi_3' = \frac{A - B}{A + B} \psi_1'}$$

- kada odredimo koeficijente A i B, zadržak je riješen!
- da li valovi bili stabilni, ne ovisi postojati c_{grm} , a to znači da δ mora biti ≥ 0 ! za totalno stabilne $\Rightarrow \delta = 0$!

- zadržak kaže $k^2 = \lambda^2 \Rightarrow c = -\frac{\beta 2k^2}{k^2 3k^2} = -\frac{2}{3} \frac{\beta}{k^2}$

$$\text{- iz uvjeta } \delta = 0 \Rightarrow \frac{\beta^2}{9k^4} - \frac{U_T^2}{3} \Rightarrow 0 \Rightarrow U_T^2 = \frac{\beta^2}{3k^4}$$

- sada dobivene iznose za U_T i ρ ubacimo u (1) i jednadžbu (1):

$$\Rightarrow (1) / ^2 \Rightarrow [(\rho - \cancel{U_{\text{kin}}})k^2 + \beta]^2 A^2 = U_T^2 k^4 B^2$$

$$\Rightarrow \left(-\frac{2}{3} \frac{\beta^2}{k^2} k^2 + \beta\right)^2 A^2 = \frac{\beta^2}{3k^4} k^4 B^2$$

$$\Rightarrow \frac{1}{9} \beta^2 A^2 = \frac{1}{3} \beta^2 B^2 \Rightarrow A^2 = 3B^2 \Rightarrow A = \pm \sqrt{3} B$$

$$\Rightarrow \psi_3' = \frac{\pm \sqrt{3} B - B}{\pm \sqrt{3} B + B} \psi_1' = \frac{\pm \sqrt{3} - 1}{\pm \sqrt{3} + 1} \psi_1' \Rightarrow \dots \Rightarrow$$

$$\Rightarrow \boxed{\psi_3' = (2 \mp \sqrt{3}) \psi_1'}$$

DZ za istu obliku moći ovimost $\omega_2'(\psi_1') = ?$

8] Iračunajte totalnu potencijalnu energiju po jedinici površine za neutralnu atmosferu ako su površinski tlak i temperatura redom $p_0 = 10^5 \text{ Pa}$ i $T_0 = 300 \text{ K}$.

$E_{\text{TOT}} = E_i + E_p$; $E_i \dots$ ovisi o temperaturi zraka
 $E_p \dots$ ovisi o težini stupca zraka
 - promatramo stupac zraka od tla do vrha atmosfere (z_A)

$$E_i = c_v \int_0^{z_A} \rho T dz ; \quad E_p = \int_0^{z_A} \rho g z dz = \int_0^{z_A} p dz = R \int_0^{z_A} \rho T dz (*)$$

- DOKAZ: $E_p = \int_0^{z_A} \rho g z dz = \int_0^{z_A} (-\frac{dp}{dz}) z dz = - \int_{p_0}^0 z dp =$
 - poravnajna integracija : $z = u \Rightarrow dz = du$
 $dp = dv \Rightarrow p = v$

$$= - \left(\frac{z p}{0, p_0} - \int_0^{z_A} p dz \right) = \int_0^{z_A} p dz \quad \checkmark \text{ OK}$$

$$\Rightarrow E_{\text{TOT}} = c_v \int_0^{z_A} \rho T dz + R \int_0^{z_A} \rho T dz = c_v \int_0^{z_A} \rho T dz + \frac{R}{c_v} c_v \int_0^{z_A} \rho T dz =$$

$$= \left(1 + \frac{R}{c_v} \right) c_v \int_0^{z_A} \rho T dz = \frac{c_v + R}{c_v} E_i = \frac{c_p}{c_v} E_i$$

$$\Rightarrow E_{\text{TOT}} = \frac{c_p}{c_v} c_v \int_0^{z_A} \rho T dz = c_p \int_0^{z_A} \rho T dz (**)$$

- za neutralnu atmosferu vrijedi $\frac{\partial \theta}{\partial z} = 0$; $\theta = T \left(\frac{p_0}{p} \right)^{\frac{R}{c_p}} / \ln / \frac{\partial}{\partial z}$
 $\rightarrow \frac{1}{\theta} \frac{\partial \theta}{\partial z} = \frac{1}{T} \frac{\partial T}{\partial z} - \frac{R}{c_p p} \left(\frac{\partial p}{\partial z} \right) = \frac{1}{T} \frac{\partial T}{\partial z} + \frac{g}{c_p} \left(\frac{\partial p}{p} \right) = \frac{1}{T} \left(\frac{\partial T}{\partial z} + \frac{g}{c_p} \right) / \theta$
 $\Rightarrow \frac{\partial \theta}{\partial z} = \frac{\theta}{T} (-\gamma + \delta) = 0 \Rightarrow \gamma = \delta$

- vidimo jednostavno linearno opadanje temperature s visinom:
 $T(z) = T_0 - \gamma z$

$$\Rightarrow T(z) = T_0 - \delta z$$

- hidrostatička: $\frac{dp}{dz} = -\rho g = -\frac{\rho}{RT} g \Rightarrow \frac{dp}{p} = -\frac{g}{RT} dz$

$$\Rightarrow \frac{dp}{p} = -\frac{g}{R} \frac{dz}{T_0 - \delta z} \quad ; \quad x = T_0 - \delta z ; \quad z=0 \Rightarrow x=T_0$$

$$\quad ; \quad dx = -\delta dz \Rightarrow dz = -\frac{dx}{\delta}$$

$$\Rightarrow \int_{p_0}^p \frac{dp}{p} = -\frac{g}{R} \int_{T_0}^{T_0 - \delta z} -\frac{dx}{\delta x} = \frac{g}{R\delta} \int_{T_0}^{T_0 - \delta z} \frac{dx}{x}$$

$$\Rightarrow \ln p / p_0 = \frac{g}{R\delta} \ln x / T_0 \Rightarrow \ln \frac{p(z)}{p_0} = \ln \left(\frac{T_0 - \delta z}{T_0} \right)^{\frac{g}{R\delta}}$$

$$\Rightarrow p(z) = p_0 \left(\frac{T_0 - \delta z}{T_0} \right)^{\frac{g}{R\delta}}$$

- vrh atmosfere: $z = z_A$; $\lim_{z \rightarrow \infty} p(z) = 0 \Rightarrow p(z_A) = 0$

$$\Rightarrow p(z_A) = 0 = p_0 \left(\frac{T_0 - \delta z_A}{T_0} \right)^{\frac{g}{R\delta}} \Rightarrow T_0 - \delta z_A = 0 \Rightarrow z_A = \frac{T_0}{\delta}$$

- sada uzmemo unos (**): $E_{\text{prot}} = C_p \int_0^{z_A} \rho T dz$, a uz unosa (*) slijedi: $\int_0^{z_A} \rho T dz = \frac{1}{R} \int_0^{z_A} p dz$

$$\Rightarrow E_{\text{prot}} = C_p \frac{1}{R} \int_0^{z_A} p dz = \frac{C_p}{R} \int_0^{z_A} p_0 \left(\frac{T_0 - \delta z}{T_0} \right)^{\frac{g}{R\delta}} dz$$

$$\Rightarrow E_{\text{prot}} = \frac{C_p p_0}{R} \int_0^{z_A} \left(1 - \frac{\delta}{T_0} z \right)^{\frac{g}{R\delta}} dz$$

$$y = 1 - \frac{\delta}{T_0} z \Rightarrow dy = -\frac{\delta}{T_0} dz \Rightarrow dz = -\frac{T_0}{\delta} dy$$

$$z=0 \Rightarrow y=1 ; \quad z=z_A \Rightarrow y = 1 - \frac{\delta}{T_0} z_A = 1 - \frac{\delta}{T_0} \frac{T_0}{\delta} = 1 - 1 = 0$$

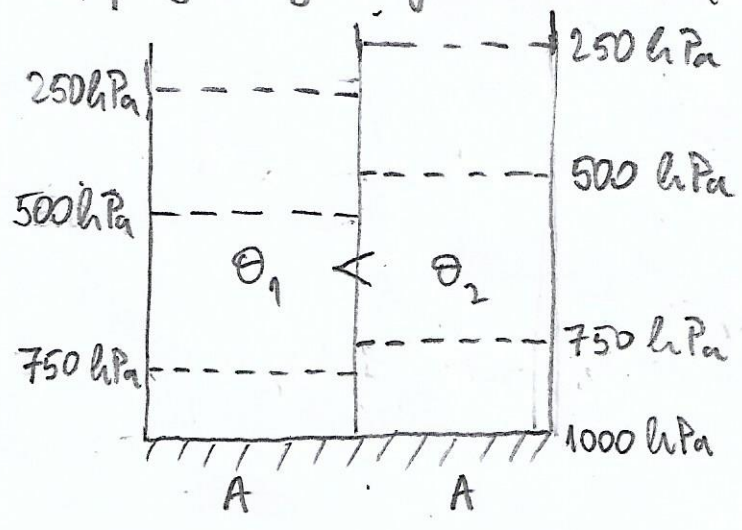
$$\Rightarrow E_{\text{prot}} = \frac{C_p p_0}{R} \int_1^0 y^{\frac{g}{R\delta}} \left(-\frac{T_0}{\delta} \right) dy = \frac{C_p p_0 T_0}{R\delta} \int_0^1 y^{\frac{g}{R\delta}} dy \quad 2,38 \cdot 10^9 \text{ J m}^{-2}$$

$$\Rightarrow E_{\text{prot}} = \frac{C_p p_0 T_0}{R\delta} \frac{1}{\frac{g}{R\delta} + 1} y^{\frac{g}{R\delta} + 1} \Big|_0^1 = \frac{C_p p_0 T_0}{R\delta} \frac{R\delta}{g + R} = \frac{p_0 T_0}{g} \frac{C_p^2}{C_p + R}$$

9] Nekoliko dvije različite vršne mase odvojene romišljenom pregradom. Svaka od njih ima konstantnu potencijalnu temperaturu $\theta_1 = 320\text{K}$ i $\theta_2 = 340\text{K}$. Također, voda vršna masa sjedi na horizontalnoj površini površine $A = 1\text{ha}$ i proteže se od tla do vrha atmosfere. Nađite iznos za dostupnu potencijalnu energiju (APE) ovog sustava.

Rj: APE ćemo dobiti kada odvojimo totalnu potencijalnu energiju sustava nakon mješovite mase od tot. pot. en. sustava s odvojenim masama: $APE = E_{p1} - E_{p2}$

1) prije mješovite zr. mase (sa pregradom) \Rightarrow topliji stupac



voda je viši od hladnijeg
- po jed. površine:

$$E_{pTOT} = \frac{\rho_0 \theta}{g} \frac{C_p^2}{C_p + R}$$

$$\theta = T \left(\frac{\rho_0}{\rho} \right)^{\frac{R}{C_p}}$$

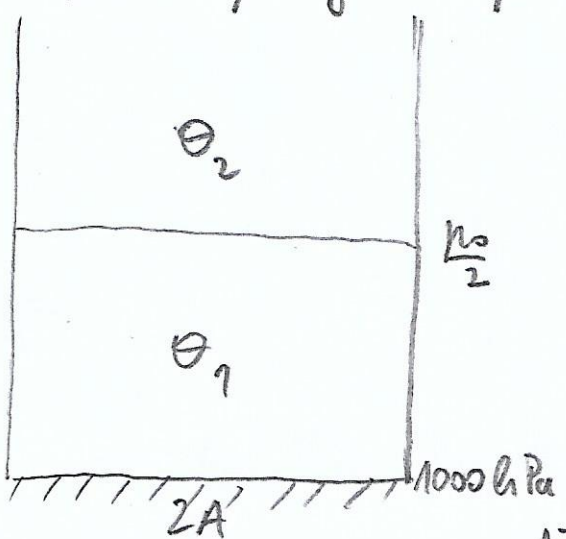
- pri tlu je $T = T_0$ i $\rho = \rho_0 \Rightarrow$

$$\Rightarrow \theta = T_0 \left(\frac{\rho_0}{\rho_0} \right)^{\frac{R}{C_p}} = T_0$$

$$\Rightarrow E_{pTOT} = \frac{\rho_0 \theta}{g} \frac{C_p^2}{C_p + R} \Rightarrow E_{pTOT A} = \frac{\rho_0 \theta}{g} \frac{C_p^2}{C_p + R} A$$

$$\Rightarrow E_{p1} = \frac{\rho_0 \theta_1}{g} \frac{C_p^2}{C_p + R} A + \frac{\rho_0 \theta_2}{g} \frac{C_p^2}{C_p + R} A = \frac{\rho_0 C_p^2 A}{g (C_p + R)} (\theta_1 + \theta_2)$$

2) pregrada se "mekine" pa će se hladniji vrti podizati ispod toplijeg. Pretpostavka je da su stupci vode identične mase pa će se granica uspostaviti točno na $\frac{h_2}{2}$.



Sada je potrebno izračunati $E_{pTOT A}$ novog sustava koji je sada stratificiran!

- općenito: $E_{\text{TOT A}} = A \frac{c_p}{c_v} E_i = A \frac{c_p}{c_v} \int_0^{z_A} \rho T dz =$
 $= -A \frac{c_p}{g} \int_{p_0}^0 T dp$ $-\frac{dp}{g}$

$$\Rightarrow E_{p2} = -2A \frac{c_p}{g} \left[\int_{p_0}^{p_0/2} \theta_1 \left(\frac{p}{p_0}\right)^{\frac{R}{c_p}} dp + \int_{p_0/2}^0 \theta_2 \left(\frac{p}{p_0}\right)^{\frac{R}{c_p}} dp \right]$$

$$\frac{p}{p_0} = x \Rightarrow dp = p_0 dx ; p_0 \rightarrow x=1 ; \frac{p_0}{2} \rightarrow x=\frac{1}{2} ; 0 \rightarrow 0$$

$$\Rightarrow E_{p2} = -2A \frac{c_p}{g} \left[\theta_1 p_0 \int_1^{\frac{1}{2}} x^{\frac{R}{c_p}} dx + \theta_2 p_0 \int_{\frac{1}{2}}^0 x^{\frac{R}{c_p}} dx \right] =$$

$$= -2A \frac{c_p}{g} \left[\theta_1 p_0 \frac{1}{\frac{R}{c_p} + 1} x^{\frac{R}{c_p} + 1} \Big|_1^{\frac{1}{2}} + \theta_2 p_0 \frac{1}{\frac{R}{c_p} + 1} x^{\frac{R}{c_p} + 1} \Big|_{\frac{1}{2}}^0 \right] =$$

$$= -2A \frac{c_p^2}{g} \frac{p_0}{R + c_p} \left[\theta_1 \left(\frac{1}{2}\right)^{\frac{R}{c_p} + 1} - \theta_1 - \theta_2 \left(\frac{1}{2}\right)^{\frac{R}{c_p} + 1} \right]$$

- kada se uveste brojke, dobije se:

$$E_{p1} \approx 5,26 \cdot 10^{13} \text{ J}$$

$$E_{p2} \approx 5,23 \cdot 10^{13} \text{ J}$$

$$\Rightarrow APE \approx 3 \cdot 10^{11} \text{ J}$$

- DZ \Rightarrow provjeriti brojke!